

# Rapid mixing of commuting spin lattice systems

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## Thermalization of quantum spin systems: From 'weak' inequalities to strong results

- **Davies thermalization** is a Markovian model for thermalization of spin lattice systems where every lattice point individually couples weakly to the heat bath at temperature  $\beta^{-1}$ .
- This dynamic drives the system to its thermal state, also called the Gibbs state:  $\rho_t \xrightarrow{t \rightarrow \infty} \sigma_\beta$ .
- **Questions:**
  1. **How fast does this happen?**  $\rightarrow t_{\text{mix}}(\epsilon) := \inf\{t \geq 0 \mid \|\rho_t - \sigma\|_1 \leq \epsilon\}$   
And how does this mixing time scale in system size  $N$ ?
  2. What is the relation to the locality structure of the system Hamiltonian/ its Gibbs state?  $\leftrightarrow$  Notions of clustering in the Gibbs state
- Consider systems on  $D$ -dim hypercubic lattices  $\Lambda \subset \mathbb{Z}^D$  with **geometrically-local, commuting**, uniformly bounded Hamiltonians:

$$\mathbb{F}_X = 0 \text{ for } \text{diam} X > r \quad \left( \begin{array}{l} H_\Lambda = \sum_{X \subset \Lambda} \Phi_X \\ [\Phi_x, \Phi_y] = 0 \quad \forall X, Y \subset \Lambda \end{array} \right. \quad (1)$$

$\rightarrow$  e.g. Toric code, ...

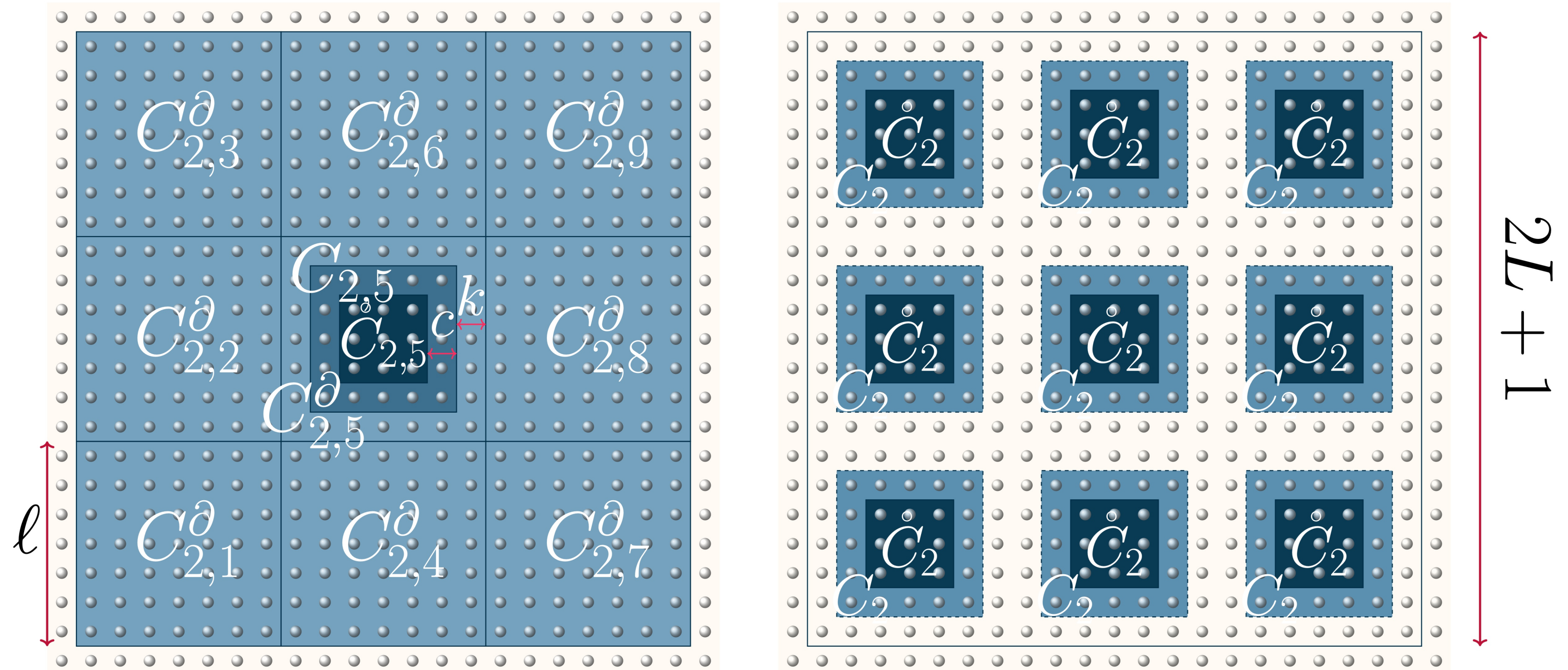


Figure 1. The first layer of a suitable covering of the  $D$ -dimensional sublattice  $[-L, L]^D$  into almost  $D - a$ -dim. regions  $\{C_{a,i}\}_{a,i}$  size  $l$  and overlap  $c$ . It is a generalization of a suitable 2D covering of [2].

### Main Results

Let  $([-L, L]^D, H_\Lambda, \beta)$  be a **commuting** and **geometrically-local** quantum spin systems on the  $D$ -dimensional lattice. Then **strong local indistinguishability** of the Gibbs state **implies** that the associated Davies dynamics

1. **Theorem 1: [Rapid thermalization]** from sLI & Gap are rapidly mixing, i.e.  $\|e^{t\mathcal{L}_\Lambda^D}(\rho) - \sigma\|_1 \leq \epsilon$  for
 
$$t \geq t_{\text{mix}}(\epsilon) = \mathcal{O}(\log(N)^D, \log \epsilon^{-1}) \quad (2)$$

when the local gap of the Davies generators  $\lambda(\mathcal{L}_X^D)$ , for  $X \subset \Lambda_L$  scales inverse polynomially with  $|X|$ .

2. **Theorem 2: [Quasi-rapid Wasserstein mixing]** from sLI satisfy quasi-rapid Wasserstein mixing, i.e.  $W_1(e^{t\mathcal{L}_\Lambda^D}(\rho), \sigma) \leq \epsilon N$  for

$$t \geq t_{W_1}(\epsilon) = \mathcal{O}(\exp(\text{poly } \log \log(N)), \log \epsilon^{-1}), \quad (3)$$

when assuming any strictly positive local gap  $\lambda(\mathcal{L}_X^D) > 0$ . We have set  $N := |\Lambda|$

*an extensive distance that quantifies local differences*

### I. Strong Local indistinguishability (sLI)

Global thermalization results rely on **lifting local rapid thermalization** to the whole lattice [3-5]. Such arguments need a notion of **spacial clustering**, i.e. stability of the Gibbs state in some region of interest  $A$  to changes of the Hamiltonian 'far away'  $C$ . To quantify this we find the notion of strong local indistinguishability very useful.

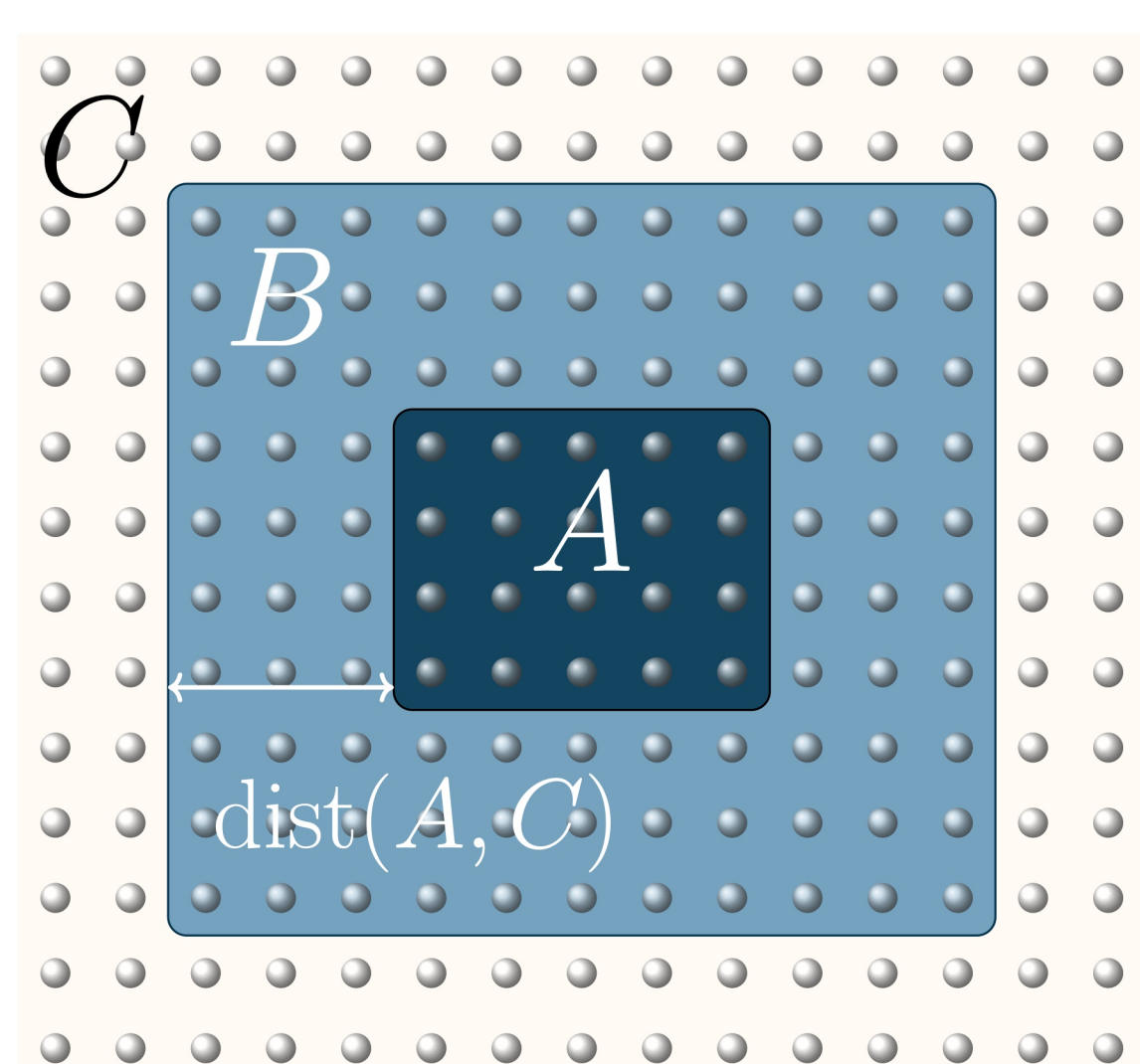


Figure 2. A region  $A$  shielded from a large region  $C$  by a ring  $B$ , as in the definition of sLI (4).

- In 1D systems of form (1) it implies rapid mixing at any temperature  $\beta > 0$  [1].

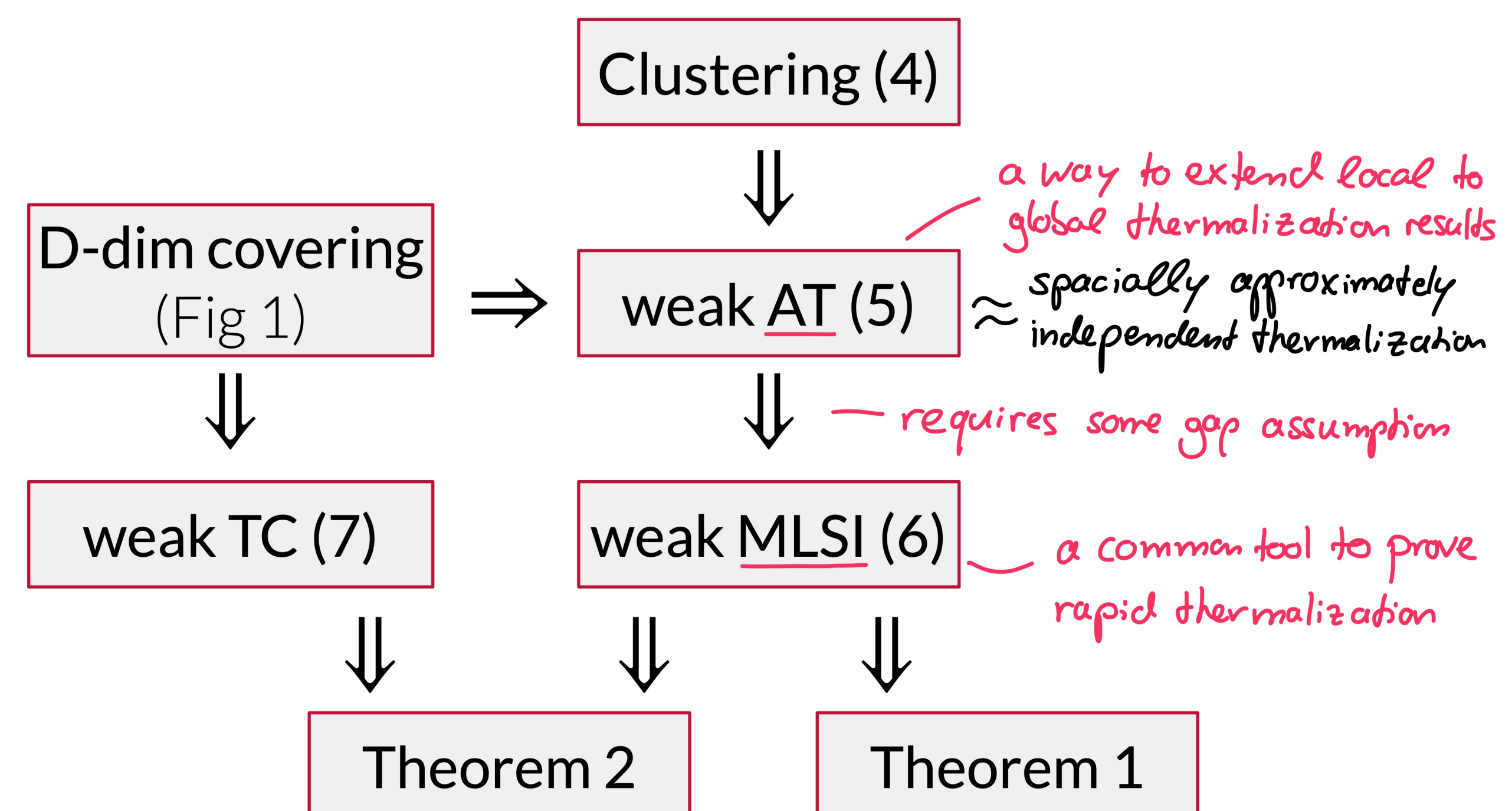
**Definition:**

$$D_{\text{max}}(\sigma_A^{ABC} \| \sigma_A^{AB}) \leq K \max\{|A|, |C|\}^\nu \exp\left(-\frac{\text{dist}(A, C)}{\xi}\right) \quad (4)$$

*exponentially decaying*

**References:** [1] Kochanowski, Alhambra, Capel, Rouzé; 2024, [2] Brandão, Kastoryano; 2019, [3] Bardet et al. 2021, [4] Capel et al. 2021, [5] Zegarliński 1994, [6] Onorati et al. 2023, [7] Chen et al. 2023 & Ding et al. 2024

### II. Proof idea and tools



- Derive a **weak approximate tensorization** statement

$$D(\rho \| \sigma^{\Lambda_L}) \leq \sum_{a,i} D(\rho \| E_{C_{a,i}}(\rho)) + \epsilon \quad (5)$$

- Derive a **weak modified logarithmic Sobolev inequality**

$$D(\rho \| \sigma^{\Lambda_L}) \leq \mathcal{O}\left((\log N \epsilon^{-1})^D\right) \text{EP}_\Lambda(\rho) + \epsilon \quad (6)$$

- Extend a **weak transport cost inequality** to  $D$ -dimensions from [6]

$$W_1(\rho, \sigma^{\Lambda_L}) \leq \mathcal{O}\left(\log\left(N \delta^{-1/2}\right)^D\right) \sqrt{ND(\rho \| \sigma)} + \mathcal{O}(\delta) \quad (7)$$

### Applications and Conclusion

- Efficient preparation of Gibbs states of geom-local commuting Hamiltonians that satisfy sLI and a polynomial gap
- Efficient estimation of local observables through  $W_1$ -decay under much weaker assumption of gap
- Efficient estimation of quantum partition functions of such systems via rapid mixing
- No go results of such systems as quantum memory devices

**Outlook:**

- Rapid mixing of non GNS-symmetric and non-commuting Lindblad generators of [7]  $\rightarrow$  efficient preparation and sampling of very general class of Gibbs states.
- Improved transport cost inequality (7) would give almost constant  $W_1$ -mixing time.